CrossMark
click for updates

Note: *In situ* measurement of vacuum window birefringence by atomic spectroscopy

Andreas Steffen, Wolfgang Alt, Maximilian Genske, Dieter Meschede,
Carsten Robens, and Andrea Alberti[a)]
Institut für Angewandte Physik, Universität Bonn, Wegelerstr. 8, D-53115 Bonn, Germany

We present an *in situ* method to measure the birefringence of a single vacuum window by means of microwave spectroscopy on an ensemble of cold atoms. Stress-induced birefringence can cause an ellipticity in the polarization of an initially linearly polarized laser beam. The amount of ellipticity can be reconstructed by measuring the differential vector light shift of an atomic hyperfine transition. Measuring the ellipticity as a function of the linear polarization angle allows us to infer the amount of birefringence Δn at the level of 10^{-8} and identify the orientation of the optical axes. The key benefit of this method is the ability to separately characterize each vacuum window, allowing the birefringence to be precisely compensated in existing vacuum apparatuses. © *2013 AIP Publishing LLC*. [http://dx.doi.org/10.1063/1.4847075]

Many experiments in quantum optics rely on an accurate control of the polarization of the laser beams.[1–6] Optical access of laser beams to ultrahigh vacuum apparatus is offered by vacuum windows, which are, in general, affected by stress-induced birefringence occurring after mounting and bake-out. While the typical values of the induced birefringence Δn are in the order of 10^{-6}, values significantly below this magnitude require special attention in mounting the vacuum viewports to avoid deformations.[7] It is, thus, important for precision applications to be able to characterize the amount of birefringence of each individual window. Knowing the amount of birefringence and the orientation of the principal axes makes it possible to avoid polarization distortions either by aligning the incoming linear polarization onto one of the optical axes, or by fully compensating the birefringence by means of optical (e.g., Soleil-Babinet compensator) or mechanical techniques.[8] However, characterizing the polarization distortion outside of the vacuum with conventional polarimeters is not sufficient to reconstruct separately the birefringence of the two vacuum viewports, which the laser beam must transit. One solution which has been proposed to obviate this problem requires employing wedged vacuum windows and picking off the beam back-reflected from the inside facet.[9] However, this is not directly applicable to standard viewports or vacuum cells.

In this note, we demonstrate an *in situ* method to reconstruct the stress-induced birefringence Δn of a vacuum window and the orientation angle θ_0 of the optical axes. Our scheme makes use of the atoms themselves as a sensitive probe to detect any ellipticity caused by mechanical stresses acting on the vacuum window, as illustrated in Figure 1(a). While varying the angle θ of the incident linear polarization, we measure the light shift δ of a hyperfine transition by means of microwave spectroscopy. We will show that the recorded signal behaves as

$$\delta \propto S_0 \sin(kL\Delta n) \sin(2(\theta - \theta_0)),\qquad(1)$$

where the proportionality constant is fully determined by the atomic properties, k is the probe laser wavevector, L is the thickness of the vacuum window, and S_0 denotes the Stokes parameter, which specifies the total intensity. Fitting the model (1) to the experimental data allows us to fully reconstruct the Jones matrix

$$M = R(\theta_0) \times \begin{pmatrix} e^{ikL\Delta n/2} & 0 \\ 0 & e^{-ikL\Delta n/2} \end{pmatrix} \times R(-\theta_0),\qquad(2)$$

which characterizes the polarization transformation caused by the vacuum window, where $R(\theta_0)$ is the two-dimensional rotation matrix with angle θ_0. By parametrizing M in the form of (2), we made the reasonable assumption that mechanical stresses cannot induce optical activity on thin windows.[10] If we represent the incident linear polarization with the Jones vector $\hat{e} = (\cos\theta, \sin\theta)$, the atoms experience the polarization $(u, v) = M\hat{e}$, which has an ellipticity $\epsilon(\theta) = 2\,\mathrm{Im}\,v^\star u = \sin(2(\theta - \theta_0))\sin(kL\Delta n)$. We assume here the conventional definition of ellipticity in terms of the Stokes parameters as $\epsilon = S_3/S_0 = (I_{\sigma+} - I_{\sigma-})/(I_{\sigma+} + I_{\sigma-})$, with I_s being the intensity of s-circularly polarized photons. Using alkali atoms as an example, we consider the differential vector light shift[11,12] caused by the probe laser beam to two hyperfine states of the ground state. We call these states $|a\rangle = |F = I + 1/2, m_F\rangle$ and $|b\rangle = |F = I + 1/2, m_F'\rangle$ in standard spectroscopic notation with the quantization axis in the direction of the probe beam. We obtain from formula (19) in Ref. 13 that the resonance frequency is shifted by

$$\delta = \alpha \frac{\nu_2 - \nu_1}{(\nu - \nu_1)(\nu - \nu_2)}(g_F' m_F' - g_F m_F)S_0\epsilon,\qquad(3)$$

where ν is the probe laser frequency, ν_1 and ν_2 denote the resonance frequencies of the D_1 and D_2 lines, g_F and g_F' represent the g-factors of the states $|a\rangle$ and $|b\rangle$, and $\alpha = c^2\Gamma_1/(32\pi^3 h\nu_1^3) \approx c^2\Gamma_2/(32\pi^3 h\nu_2^3)$ is a proportionality constant depending on the atomic parameters such as the natural decay rates Γ_1 and Γ_2 of the doublet states, as well as the speed of light c and Planck constant h.[14] Inserting the expression of the ellipticity $\epsilon(\theta)$ in (3), we readily obtain the
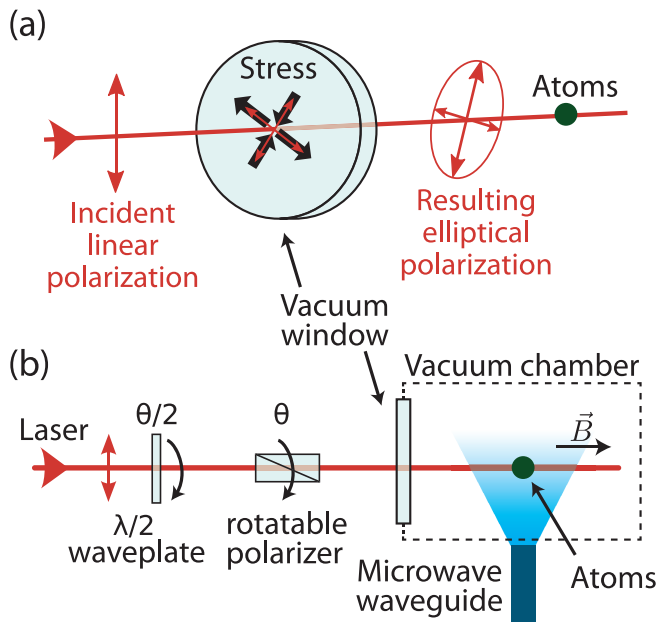

FIG. 1. (a) Polarization distortion caused by stress-induced birefringence in vacuum windows. A linear polarization entering at an angle to the principal stress axes acquires a degree of ellipticity, which in turn induces a detectable shift of a hyperfine transition of the atoms. (b) Measurement scheme. A half-wave plate and rotatable polarizer allow the linear-polarization of a probe laser beam to be prepared at any arbitrary angle θ. While atoms are homogeneously irradiated by the probe laser, the resonance frequency of the hyperfine transition is measured as a function of θ using microwave spectroscopy. The magnetic field $\vec{B}$ defines the quantization axis.
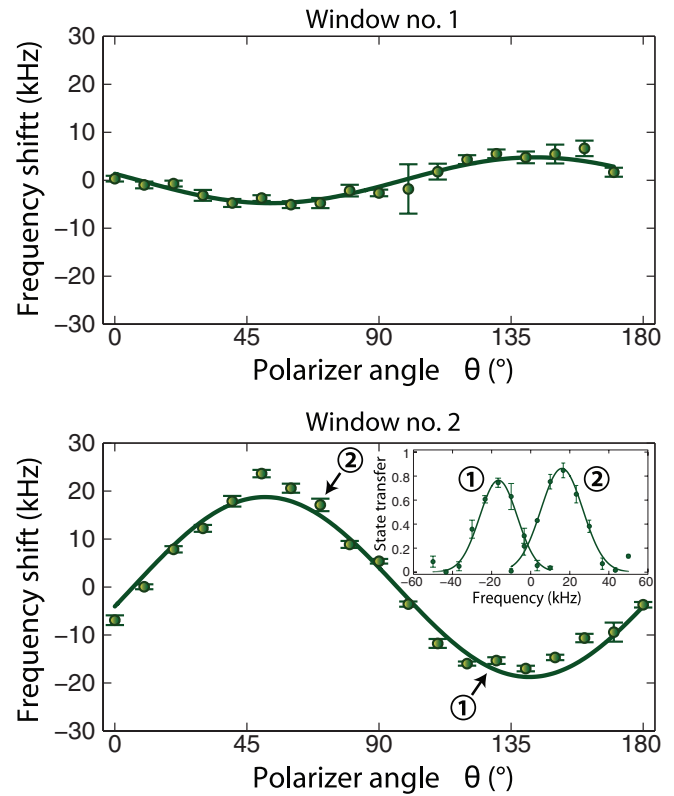

FIG. 2. Measured differential light shift as a function of the polarization angle θ. The transition frequency varies in a sinusoidal way according to (1). Shown are the data with 1 sigma error bars and a sinusoidal fit. The offset of 9.199810 GHz, including the contribution from differential scalar light shift and magnetic fields, has been subtracted. The origin $\theta = 0$ is arbitrarily set by the polarizer's orientation in its mount. Inset: Two measured spectra, corresponding to the two indicated data points.

expression in (1). Formula (3) prompts a few considerations about the sensitivity $\beta = \delta/(\epsilon S_0)$ of atoms to elliptical polarization per unit of laser intensity: The sensitivity is proportional to the line doublet splitting $\nu_2 - \nu_1$ (in general, spin-orbit coupling), which makes heavier atoms more sensitive than lighter atoms. In addition, probe frequencies ν close to the optical resonance are more favorable, since the sensitivity scales approximately with $1/(\nu - \nu_0)^2$ for large detunings. But more importantly, a pair of magnetic-field sensitive states $|a\rangle$ and $|b\rangle$ must be chosen to measure the vector light shift, so that $g'_F m'_F \neq g_F m_F$; this excludes, for instance, employing the so-called clock states.

In our experiment, a small ensemble of ^{133}Cs atoms is cooled in a magneto-optical trap and subsequently transferred to an optical dipole trap. The trap stems from the probe laser beam itself, which has a wavelength of 866 nm and is tightly focused on the atoms. The light polarization is precisely set by a rotatable Glan-laser polarizer ($0.5°$ angle reproducibility) with a preceding half-wave plate used to maximize the transmission through the polarizer, as shown in Figure 1(b); the purity of the incident polarization is determined by the polarizer's extinction ratio of 10^{-6}. We tested two different vacuum windows: a cuboid glass cell made of Corning Vycor® 7913 (manufactured by Hellma) and a standard CF63 viewport of made of Spectrosil® 2000. We will call them, for convenience, window no. 1 and no. 2. Both are 5 mm thick, while the probe intensity experienced by the atoms during the microwave spectroscopy is ∼12 W/mm^2 for window no. 1 and ∼ 6 W/mm^2 for window no. 2. A magnetic field of about 3 G

is added along the probe beam's direction to define the quantization axis, and the atoms are prepared by optical pumping to the state $|b\rangle = |F = 4, m_F = 4\rangle$. We drive the transition to the state $|a\rangle = |F = 3, m_F = 3\rangle$ using microwave pulses at varying frequencies around 9.2 GHz. With this choice of the Zeeman states, we obtain the largest sensitivity to elliptical polarization (as $g'_F = -g_F$), resulting in $\beta = -50$ kHz/(W/mm^2) for our case. The recorded Fourier-limited spectra exhibit a FWHM of about 30 kHz.[15]

We obtain the light shift δ from the recorded spectra as the displacement of the center frequency for different angles θ, see Figure 2. The data exhibit the expected sinusoidal behavior $\delta = \delta_0 \sin(2(\theta - \theta_0))$ with the amplitude specified by $\delta_0 = \beta S_0 \sin(kL\Delta n)$. The zero-crossing determines the angle θ_0 of one of the optical axes, while the inferred value of δ_0 yields the birefringence Δn. The results are listed in Table I, showing that the birefringence magnitude and the orientation of optical axes can be extracted with good precision. We attain in our case a precision level of 10^{-8} even though our apparatus can only probe small ensembles of less than 30 atoms at each iteration 2 s cycle time. Because the light shift δ_0 is proportional to S_0, we must take into account a conservative 10% systematic uncertainty arising from the probe beam's intensity, which depends on the precise position and geometry of the beam waist. If required, S_0 can be precisely calibrated by introducing in the incident

TABLE I. Measurement results with statistical errors. From top to bottom, the maximum transition shift (for $\theta = \theta_0 + 45°$) and the zero-crossing angle are reported. Inferred from δ_0 are the maximal ellipticity and the magnitude of birefringence, with the systematic error coming from the uncertainty on the total light intensity.

Quantity		Window no. 1			Window no. 2		
		Value	Statistical	Systematic	Value	Statistical	Systematic
δ_0 (kHz)		4.8	± 0.4	–	18.8	± 0.9	–
θ_0 (°)		98	± 2	–	6	± 1	–
max ϵ	(10^{-3})	8.2	± 0.7	± 0.9	64	± 3	± 6
Δn	(10^{-7})	2.3	± 0.2	± 0.2	17.7	± 0.8	± 1.8

polarization a known amount of ellipticity (up to $\epsilon = 1$) and measuring the resulting light shift. From our measurements, we learn that the glass cell (window no. 1) exhibits a smaller birefringence than the commercial viewport (window no. 2). A possible explanation is that the frontal surface of the cell is further away from the mounting flange, where most of mechanical stresses are localized.[7] If we treat the window and its internal stresses as planar and assume the stress distribution to be homogeneous over the size of the laser beam, the stress distribution can be expressed in terms of two orthogonal principal stresses Σ_1, Σ_2, as illustrated in Figure 1(a). We can hence use the stress optic law to derive the differential stress $\Sigma_D = \Sigma_1 - \Sigma_2 = \Delta n / C$ from the measured birefringence Δn, with C being the stress optic coefficient of the material.[16] In case of fused silica, which is the main constituent of our windows, C is around 3.5×10^{-12} Pa^{-1}, yielding a Σ_D of 70 kPa for window no. 1 and 500 kPa for window no. 2. These are typical values occurring in a vacuum cell and a commercial Conflat viewport, respectively.[8] In critical applications demanding ultralow stress-induced birefringence, a special glass like Schott SF57 could provide two orders of magnitude suppression due to its unusually small stress-optic coefficient of 2×10^{-14} Pa^{-1}.[17]

In conclusion, we have presented a novel *in situ* method to measure with high precision the ellipticity caused by the stress-induced birefringence. This method allows us to align the incident linear polarization along one of the optical axes,

virtually canceling the effect of birefringence. Alternatively, one could employ a waveplate with the retardance set to $kL\Delta n$ or a Soleil-Babinet compensator in order to preserve any arbitrary incident polarization.

[1] C. Monroe, D. M. Meekhof, B. E. King, and D. J. Wineland, Science **272**, 1131 (1996).

[2] O. Mandel, M. Greiner, A. Widera, T. Rom, T. W. Hänsch, and I. Bloch, Nature (London) **425**, 937 (2003).

[3] P. J. Lee, M. Anderlini, B. L. Brown, J. Sebby-Strabley, W. D. Phillips, and J. V. Porto, Phys. Rev. Lett. **99**, 020402 (2007).

[4] J.-M. Choi and D. Cho, J. Phys.: Conf. Ser. **80**, 012037 (2007).

[5] R. Chicireanu, K. D. Nelson, S. Olmschenk, N. Lundblad, A. Derevianko, and J. V. Porto, Phys. Rev. Lett. **106**, 063002 (2011).

[6] A. Steffen, A. Alberti, W. Alt, N. Belmechri, S. Hild, M. Karski, A. Widera, and D. Meschede, Proc. Natl. Acad. Sci. U.S.A. **109**, 9770 (2012).

[7] N. Solmeyer, K. Zhu, and D. S. Weiss, Rev. Sci. Instrum. **82**, 066105 (2011).

[8] A. A. Studna, D. E. Aspnes, L. T. Florez, B. J. Wilkens, J. P. Harbison, and R. E. Ryan, J. Vac. Sci. Technol. A **7**, 3291 (1989).

[9] B. Park, A. Sushkov, and D. Budker, Rev. Sci. Instrum. **79**, 013108 (2008).

[10] M represents a SU(2) rotation on the Poincaré sphere by an angle $kL\Delta n$ around the axis $(-\sin 2\theta_0, \cos 2\theta_0, 0)$, which excludes optical activity.

[11] I. Deutsch and P. Jessen, Phys. Rev. A **57**, 1972 (1998).

[12] J. M. Geremia, J. K. Stockton, and H. Mabuchi, Phys. Rev. A **73**, 042112 (2006).

[13] R. Grimm, M. Weidemüller, and Y. Ovchinnikov, Adv. At., Mol., Opt. Phys. **42**, 95 (2000).

[14] For simplicity's sake, formula (3) assumes that the hyperfine interaction is not resolved for sufficiently large detunings. In this approximation, the vector polarizability operator is proportional to the electronic spin[11] and, because $L = 0$ in the ground state and because of the rotational symmetry along the quantization axis, it can be expressed through $g_F m_F$ using the Landé projection theorem. Here, g_F does not include the nuclear contribution.

[15] D. Meschede and A. Rauschenbeutel, Adv. At., Mol., Opt. Phys. **53**, 75 (2006).

[16] M. Born and E. Wolf, *Principles of Optics: Electromagnetic Theory of Propagation, Interference and Diffraction of Light* (Cambridge University Press, 1999).

[17] T. McMillan, P. Taborek, and J. E. Rutledge, Rev. Sci. Instrum. **75**, 5005 (2004).